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Evidence for $e^+e^- \rightarrow \gamma\chi_{c1,2}$ at center-of-mass energies from 4.009 to 4.360 GeV*

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Abstract: Using data samples collected at center-of-mass energies of $\sqrt{s}=4.009, 4.230, 4.260$, and 4.360 GeV with the BESIII detector operating at the BEPCII collider, we perform a search for the process $e^+e^- \rightarrow \gamma\chi_{cJ}$ ($J=0, 1, 2$)

and find evidence for $e^+e^- \rightarrow \gamma\chi_{c1}$ and $e^+e^- \rightarrow \gamma\chi_{c2}$ with statistical significances of 3.0σ and 3.4σ , respectively. The Born cross sections $\sigma^B(e^+e^- \rightarrow \gamma\chi_{cJ})$, as well as their upper limits at the 90% confidence level (C.L.) are determined at each center-of-mass energy.

Key words: heavy quarkonia, decays of hadronic, electron-positron collisions, hadron production by

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1 Introduction

The charmonium-like state $Y(4260)$ was first observed in the initial state radiation (ISR) process $e^+e^- \rightarrow \gamma_{\text{ISR}}\pi^+\pi^-J/\psi$ by BaBar [1], and later confirmed by the CLEO [2] and Belle [3] experiments. Recently, both BaBar and Belle updated their results with full data sets, and further confirmed the existence of the $Y(4260)$ [4, 5]. Since it is produced through ISR in e^+e^- annihilation, the $Y(4260)$ has the quantum numbers $J^{PC}=1^{--}$. However, there seems to be no $c\bar{c}$ slot available for the $Y(4260)$ in the conventional charmonium family [6]. In addition, a number of unusual features, such as a strong coupling to hidden-charm final states, suggest that the $Y(4260)$ is a non-conventional $c\bar{c}$ meson. Possible interpretations of this state can be found in Refs. [7–11], but all need further experimental input.

Most of the previous studies of the $Y(4260)$ have utilized hadronic transitions. Except for the clear signal observed in the $\pi^+\pi^-J/\psi$ decay mode, the Belle experiment failed to find evidence of the $Y(4260)$ via the $e^+e^- \rightarrow \gamma_{\text{ISR}}\eta J/\psi$ process [12]. Based on 13.2 pb^{-1} of e^+e^- data collected at $\sqrt{s}=4.260 \text{ GeV}$, the CLEO experiment investigated fourteen hadronic decay channels, but the only charmonium channels with more than 3σ statistical significance are $\pi^+\pi^-J/\psi$, $\pi^0\pi^0J/\psi$ and K^+K^-J/ψ [13]. The BESIII Collaboration first observed the process $e^+e^- \rightarrow \gamma X(3872)$ using data samples taken between $\sqrt{s}=4.009$ and 4.420 GeV [14], which strongly supports the existence of the radiative transition decays of the $Y(4260)$. To further understand the nature of the $Y(4260)$ state, an investigation into the radiative transitions between the $Y(4260)$ and other lower mass charmonium states, like the χ_{cJ} ($J=0, 1, 2$), is important [15, 16]. The cross sections of $e^+e^- \rightarrow \gamma\chi_{cJ}$ have been evaluated theoretically within the framework of non-relativistic quantum chromodynamics (NRQCD) [16]. Experimentally, the only existing investigation comes from the CLEO experiment [13], which did not observe a signal. The large data sample collected with the BESIII detector provides a good opportunity to deeply investigate these decay modes, which may shed more light on the properties of the $Y(4260)$.

In this paper, we report on a search for $e^+e^- \rightarrow \gamma\chi_{cJ}$ ($J=0, 1, 2$) based on the large e^+e^- annihilation data samples collected with the BESIII detector at center-of-mass energies (CME) $\sqrt{s}=4.009, 4.230, 4.260$, and

4.360 GeV , where the χ_{cJ} is reconstructed by its $\gamma J/\psi$ decay mode, and the J/ψ is by its decay to $\mu^+\mu^-$. The decay $J/\psi \rightarrow e^+e^-$ is not considered in this analysis due to the high background of Bhabha events. The corresponding luminosities of the data samples at different CME used in this analysis are listed in Table 1.

Table 1. The center-of-mass energy and Luminosity of each data sample.

\sqrt{s}/GeV	luminosity/ pb^{-1}
4.009	482
4.230	1047
4.260	826
4.360	540

2 BESIII detector and Monte Carlo

The BESIII detector at the BEPCII collider [17] is a large solid-angle magnetic spectrometer with a geometrical acceptance of 93% of 4π solid angle consisting of four main components. The innermost is a small-cell, helium-based (40% He, 60% C_3H_8) main drift chamber (MDC) with 43 layers providing an average single-hit resolution of $135 \mu\text{m}$. The resulting charged-particle momentum resolution for a 1 T magnetic field setting is 0.5% at $1.0 \text{ GeV}/c$, and the resolution on the ionization energy loss information (dE/dx) is better than 6%. The next detector, moving radially outwards, is a time-of-flight (TOF) system constructed of 5 cm thick plastic scintillators, with 176 detectors of 2.4 m length in two layers in the barrel and 96 fan-shaped detectors in the end-caps. The barrel (end-cap) time resolution of 80 ps (110 ps) provides a 2σ K/π separation for momenta up to $1.0 \text{ GeV}/c$. Continuing outward, we have an electromagnetic calorimeter (EMC) consisting of 6240 CsI(Tl) crystals in a cylindrical barrel structure and two end-caps. The energy resolution at 1.0 GeV is 2.5% (5%) and the position resolution is 6 mm (9 mm) in the barrel (end-caps). Finally, the muon counter consists of 1000 m^2 of Resistive Plate Chambers in nine barrel and eight end-cap layers, which provides a 2 cm position resolution.

A GEANT4 [18] based Monte Carlo (MC) simulation software, which includes the geometric description of the detector and the detector response, is used to optimize the event selection criteria, determine the detection efficiency, and estimate the potential backgrounds. Signal MC samples of $e^+e^- \rightarrow \gamma\chi_{cJ}$ are generated for

each CME according to the electric-dipole (E1) transition assumption [19]. Effects of ISR are simulated with KKMC [20] by assuming that $\gamma\chi_{cJ}$ is produced via $Y(4260)$ decays, where the $Y(4260)$ is described by a Breit-Wigner function with resonance parameters from the world average [21]. For the background studies, an ‘inclusive’ $Y(4260)$ MC sample equivalent to 500 pb^{-1} integrated luminosity is generated which includes the $Y(4260)$ resonance, ISR production of the known vector charmonium states, and events driven by QED processes. The known decay modes are generated with EvtGen [19] with branching fractions set to their world average values in the Particle Data Group (PDG) [22], and the remaining events are generated with Lundcharm [23] or PYTHIA [24].

3 Event selection

Charged tracks are reconstructed in the MDC. For each good charged track, the polar angle must satisfy $|\cos\theta| < 0.93$, and the point of closest approach to the interaction point must be within $\pm 10 \text{ cm}$ in the beam direction and within $\pm 1 \text{ cm}$ in the plane perpendicular to the beam direction. The number of good charged tracks is required to be two with a zero net charge. Charged tracks are identified as muons if they have $E/p < 0.35$ and $p > 1.0 \text{ GeV}/c$, where E is the energy deposited in the EMC and p is the momentum measured by the MDC.

Photons are reconstructed from isolated showers in the EMC that are at least 20 degrees away from any of the charged tracks. To improve the reconstruction efficiency and the energy resolution, the energy deposited in the nearby TOF counters is included. Photon candidates are required to have energy greater than 25 MeV in the EMC barrel region ($|\cos\theta| < 0.8$), and 50 MeV in the end-cap region ($0.86 < |\cos\theta| < 0.92$). In order to suppress electronic noise and energy deposits that are unrelated to the event, the EMC time t of the photon candidates must be in coincidence with collision events within the range $0 \leq t \leq 700 \text{ ns}$. At least two photon candidates in the final state are required.

To improve the momentum resolution and to reduce backgrounds, a kinematic fit with five constraints (5C-fit) is performed under the $e^+e^- \rightarrow \gamma\gamma\mu^+\mu^-$ hypothesis, imposing overall energy and momentum conservation and constraining the invariant mass of $\mu^+\mu^-$ to the nominal J/ψ mass. Candidates with a $\chi^2_{5C} < 40$ are selected for further analysis. If more than one candidate occurs in an event, the one with the smallest χ^2_{5C} is selected. Due to the kinematics of the reaction, the first radiative photon from $e^+e^- \rightarrow \gamma\chi_{cJ}$ has a high energy while the second radiative photon from $\chi_{cJ} \rightarrow \gamma J/\psi$ has a lower energy at $\sqrt{s}=4.230, 4.260$, and 4.360 GeV . The invariant mass of

the low energy photon and J/ψ , $M_{\gamma J/\psi}$, is used to search for χ_{cJ} signals. However, for the data sample taken at $\sqrt{s}=4.009 \text{ GeV}$, there is an overlap of the energy distributions of the photons from $e^+e^- \rightarrow \gamma\chi_{c1,2}$ and from $\chi_{c1,2}$ decays, as shown in Fig. 1. To separate the overlapping photon spectra, the energy of photons from $\chi_{c1,2}$ decays is further required to be less than 0.403 GeV at $\sqrt{s}=4.009 \text{ GeV}$.

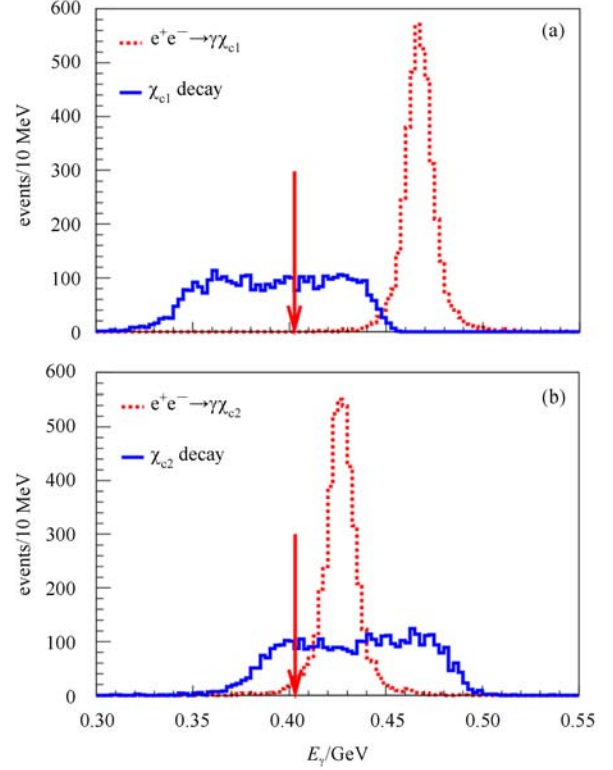


Fig. 1. The distributions of photon energies in the laboratory frame from $e^+e^- \rightarrow \gamma\chi_{c1,2}$ and from $\chi_{c1,2}$ decays in the exclusive MC samples of $e^+e^- \rightarrow \gamma\chi_{c1}$, $\chi_{c1} \rightarrow \gamma J/\psi$ (a) and $e^+e^- \rightarrow \gamma\chi_{c2}$, $\chi_{c2} \rightarrow \gamma J/\psi$ (b) at $\sqrt{s}=4.009 \text{ GeV}$. Dashed lines stand for the first radiative photons from $e^+e^- \rightarrow \gamma\chi_{c1,2}$ and solid lines for the second radiative photons from $\chi_{c1,2}$ decays.

4 Background study

The potential backgrounds from $e^+e^- \rightarrow P + J/\psi$, $P \rightarrow \gamma\gamma$ ($P = \pi^0, \eta$, or η') can be rejected by requiring $|M_{\gamma\gamma} - M_{\pi^0}| > 0.025 \text{ GeV}/c^2$, $|M_{\gamma\gamma} - M_{\eta}| > 0.03 \text{ GeV}/c^2$ and $|M_{\gamma\gamma} - M_{\eta'}| > 0.02 \text{ GeV}/c^2$, where $M_{\gamma\gamma}$ is the invariant mass of two selected photons. The background from $e^+e^- \rightarrow \gamma_{\text{ISR}}\psi(2S)$, $\psi(2S) \rightarrow \gamma\chi_{cJ}$ is rejected by applying the 5C kinematic fit. After imposing all the selection criteria above, the remaining dominant background is from radiative dimuon events, which is not expected to peak in the $M_{\gamma J/\psi}$ distribution. This has been validated by a dedicated simulation study. For other remaining

backgrounds, such as $e^+e^- \rightarrow \pi^0\pi^0 J/\psi$, only 3.8 events (normalized to data luminosity) survive and can be neglected.

5 Fit to the mass spectrum

The resulting $M_{\gamma J/\psi}$ distributions, after applying the above selection criteria, at $\sqrt{s}=4.009, 4.230, 4.260$ and 4.360 GeV are shown in Fig. 2. An unbinned maximum likelihood fit of the $M_{\gamma J/\psi}$ distribution is performed to

extract the numbers of χ_{cJ} signal events. In the fit, the shapes of the χ_{cJ} signals are described by double Gaussian functions, where the means and the standard deviations of the double Gaussian functions are determined from a fit to the corresponding signal MC samples at $\sqrt{s}=4.260$ GeV. These shapes are also used for the other three CME points, as the resolution varies only mildly between $\sqrt{s}=4.009$ – 4.360 GeV. This has been validated by MC simulation. Since the dominant background

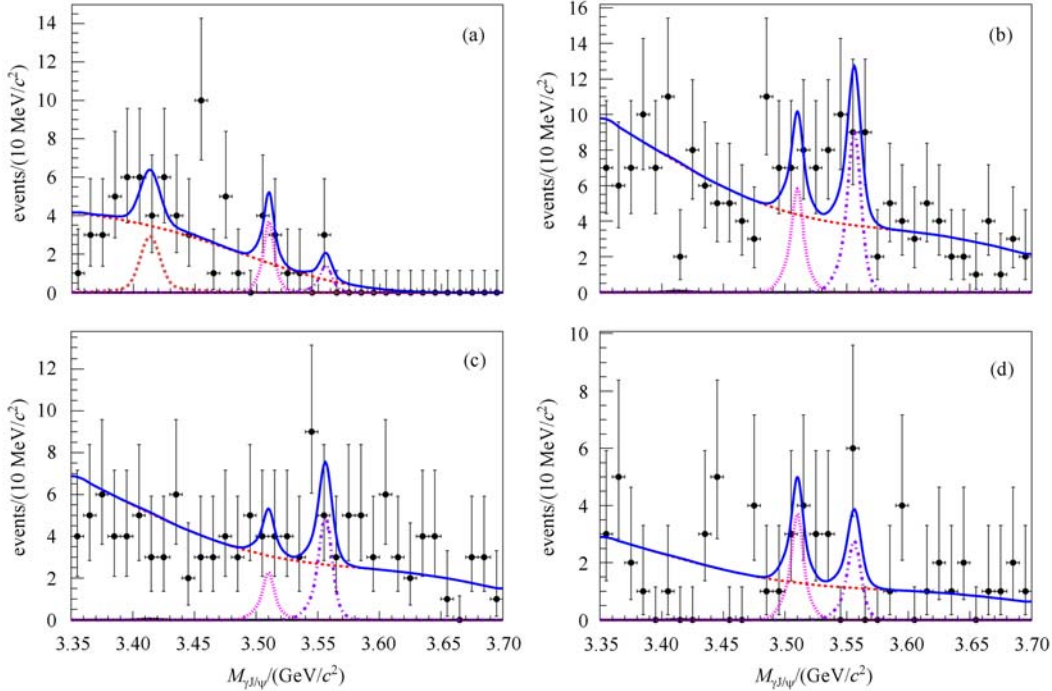


Fig. 2. The distribution of $\gamma J/\psi$ invariant mass, $M_{\gamma J/\psi}$, and fit results for data at $\sqrt{s}=4.009$ (a), 4.230 (b), 4.260 (c) and 4.360 GeV (d). The solid lines show the total fit results. The χ_{cJ} signals are shown as dashed lines, dotted lines, and dash-dotted lines, for $J = 0, 1$, and 2 , respectively. The backgrounds are indicated by red dashed lines.

Table 2. The results on $e^+e^- \rightarrow \gamma\chi_{cJ}$ Born cross section measurement. Shown in the table are the significance σ , detection efficiency ϵ , number of signal events from the fits N^{obs} , radiative correction factor $(1+\delta^r)$, vacuum polarization factor $(1+\delta^v)$, upper limit (at the 90% confidence level (C.L.)) on the number of signal events N^{UP} , Born cross section σ^{B} and upper limit (at the 90% C.L.) on the Born cross section σ^{UP} at different CME points. The first uncertainty of the Born cross section is statistical, and the second systematic.

\sqrt{s}/GeV		N^{obs}	significance (σ)	N^{UP}	ϵ (%)	$1+\delta^r$	$1+\delta^v$	$\sigma^{\text{UP}}/\text{pb}$	$\sigma^{\text{B}}/\text{pb}$
4.009	χ_{c0}	7.0 ± 6.6	1.6	18	36.4 ± 0.2			179	$65.0 \pm 61.3 \pm 4.2$
	χ_{c1}	4.4 ± 2.6	2.2	9	23.4 ± 0.1	0.738	1.04	5.3	$2.4 \pm 1.4 \pm 0.2$
	χ_{c2}	1.8 ± 1.7	1.5	6	8.7 ± 0.1			18	$4.7 \pm 4.4 \pm 0.6$
4.230	χ_{c0}	0.2 ± 2.3	0.0	7	37.2 ± 0.2			26	$0.7 \pm 8.0 \pm 0.1$
	χ_{c1}	6.7 ± 4.3	1.9	14	44.4 ± 0.2	0.840	1.06	1.7	$0.7 \pm 0.5 \pm 0.1$
	χ_{c2}	13.3 ± 5.2	2.9	22	42.0 ± 0.2			5.0	$2.7 \pm 1.1 \pm 0.3$
4.260	χ_{c0}	0.1 ± 1.9	0.0	5	36.7 ± 0.2			25	$0.5 \pm 8.8 \pm 0.1$
	χ_{c1}	3.0 ± 3.0	1.1	7	42.7 ± 0.2	0.842	1.06	1.1	$0.4 \pm 0.4 \pm 0.1$
	χ_{c2}	7.5 ± 3.9	2.3	14	41.7 ± 0.2			4.2	$2.0 \pm 1.1 \pm 0.2$
4.360	χ_{c0}	0.1 ± 0.7	0.0	3	32.4 ± 0.2			23	$0.7 \pm 5.0 \pm 0.1$
	χ_{c1}	5.2 ± 4.9	2.4	10	31.7 ± 0.2	0.943	1.05	2.9	$1.4 \pm 1.3 \pm 0.1$
	χ_{c2}	4.4 ± 4.5	2.0	9	30.3 ± 0.2			5.0	$2.3 \pm 2.3 \pm 0.2$

Table 3. Summary of systematic uncertainties at $\sqrt{s}=4.009, 4.230, 4.260$, and 4.360 GeV(%).

\sqrt{s}/GeV	4.009			4.230			4.260			4.360		
sources	χ_{c0}	χ_{c1}	χ_{c2}	χ_{c0}	χ_{c1}	χ_{c2}	χ_{c0}	χ_{c1}	χ_{c2}	χ_{c0}	χ_{c1}	χ_{c2}
luminosity	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0
tracking efficiency	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0
photon detection	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0	2.0
kinematic fit	0.6	0.6	0.6	0.6	0.6	0.6	0.6	0.6	0.6	0.6	0.6	0.6
branching ratio	4.8	3.6	3.7	4.8	3.6	3.7	4.8	3.6	3.7	4.8	3.6	3.7
vacuum polarization factor	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5
χ_{cJ} mass resolution	0.3	2.0	7.4	0.0	7.7	7.8	0.0	4.3	6.5	0.0	1.1	2.0
χ_{cJ} mass	0.0	0.9	1.4	0.0	0.3	0.2	0.0	0.1	0.1	0.0	0.3	0.4
MC model	0.0	2.2	3.9	0.0	1.2	3.3	0.0	1.9	2.9	0.0	1.5	2.1
fit range	0.1	2.2	2.6	0.0	1.5	2.3	0.0	3.1	2.5	0.0	3.0	3.7
background shape	0.0	3.1	5.6	0.0	0.7	0.3	0.0	1.1	0.4	0.0	0.9	0.1
radiative correction factor	3.0	2.7	3.6	2.6	3.1	2.1	3.5	2.1	2.5	1.8	3.4	3.8
total	6.5	7.3	12.1	6.3	9.8	10.2	6.7	7.7	9.3	6.0	6.9	7.7

comes from radiative dimuon events, the corresponding MC simulation is used to represent the background shape. To reduce the effect of statistical fluctuations, the dimuon MC shape is smoothed before it is taken as the background function. Fig. 2 also shows the fitted results for the $M_{\gamma J/\psi}$ distribution at different CME. The number of fitted χ_{cJ} signal events, as well as the corresponding statistical significances (calculated by comparing the fit log likelihood values with and without the χ_{cJ} signal) at the four CME points are listed in Table 2. The same fit is applied to the sum of $M_{\gamma J/\psi}$ distributions of the four CME points. The statistical significances for χ_{c0} , χ_{c1} and χ_{c2} are found to be 1.2σ , 3.0σ and 3.4σ , respectively. The goodness of the fit is estimated by using a χ^2 test method with the data distributions regrouped to ensure that each bin contains more than 7 events. The test gives $\chi^2/\text{d.o.f}=39.7/32$, where d.o.f is the number of degrees of freedom. As a test, we perform similar analyses to control samples from the J/ψ sideband regions, $2.917 < M_{\mu^+\mu^-} < 3.057$ GeV/ c^2 and $3.137 < M_{\mu^+\mu^-} < 3.277$ GeV/ c^2 , by constraining the invariant mass of $\mu^+\mu^-$ to 2.987 or 3.207 GeV/ c^2 in 5C-fit, and find no obvious χ_{cJ} signals.

6 Results

The Born cross section at different CME is calculated with

$$\sigma^B(e^+e^- \rightarrow \gamma\chi_{cJ}) = \frac{N^{\text{obs}}}{\mathcal{L} \cdot (1+\delta^r) \cdot (1+\delta^v) \cdot \mathcal{B} \cdot \epsilon}, \quad (1)$$

where N^{obs} is the number of observed events obtained from the fit, \mathcal{L} is the integrated luminosity, $1+\delta^r$ is the radiative correction factor for χ_{cJ} with the assumption that the $e^+e^- \rightarrow \gamma\chi_{cJ}$ cross section follows the Y(4260) Breit-Wigner line shape [25], $1+\delta^v$ is the vacuum polarization factor [26], \mathcal{B} is the combined branching ratio of $\chi_{cJ} \rightarrow \gamma J/\psi$ and $J/\psi \rightarrow \mu^+\mu^-$, and ϵ is the detection

efficiency. The detection efficiencies, radiative correction factors as well as the calculated Born cross sections at different CME are shown in Table 2. The much lower efficiencies for $\chi_{c1,2}$ at $\sqrt{s}=4.009$ GeV are due to the requirement on the photon energy used to separate the overlapping photon spectra as described in Section 3.

Since the χ_{cJ} signals are not statistically significant at the individual CME points, we also give in Table 2 the upper limits on the Born cross sections at the 90% confidence level (C.L.) under the assumption that no signals are present. The upper limits are derived using a Bayesian method [21], where the efficiencies are lowered by a factor of $(1-\sigma_{\text{sys}})$ to take systematic uncertainties into account.

We also perform a simultaneous fit to the $M_{\gamma J/\psi}$ distribution at $\sqrt{s}=4.009, 4.230, 4.260$, and 4.360 GeV, assuming the production cross section of $e^+e^- \rightarrow \gamma\chi_{cJ}$ at a different CME point follows the line shape of the Y(4260) state. In the fit, the line shapes of the χ_{cJ} signals and the background are the same as those in the previous fits, and the number of χ_{cJ} events at each CME point is expressed as a function of $\epsilon_{\text{c.m.}} \mathcal{L}_{\text{c.m.}} R_{\text{c.m.}} (1+\delta^r)$, where $\epsilon_{\text{c.m.}}$ and $\mathcal{L}_{\text{c.m.}}$ are the detection efficiency and luminosity, respectively, and $R_{\text{c.m.}}$ is the ratio of the cross section calculated with the Y(4260) line shape (a Breit-Wigner function with parameters fixed to the PDG values) at different CME points to that at $\sqrt{s}=4.260$ GeV. The corresponding fit result is shown in Fig. 3. The goodness of the fit is $\chi^2/\text{d.o.f}=53.3/40$ and the statistical significances for χ_{c0} , χ_{c1} and χ_{c2} signals are 0σ , 2.4σ and 4.0σ , respectively. We also found that $\Gamma_{e^+e^-} \times \mathcal{B}(Y(4260) \rightarrow \gamma\chi_{c1}) = (0.11 \pm 0.06)$ eV and $\Gamma_{e^+e^-} \times \mathcal{B}(Y(4260) \rightarrow \gamma\chi_{c2}) = (0.33 \pm 0.11)$ eV.

7 Systematic uncertainties

The systematic uncertainties in the cross section measurements of $e^+e^- \rightarrow \gamma\chi_{cJ}$ are caused by various sources

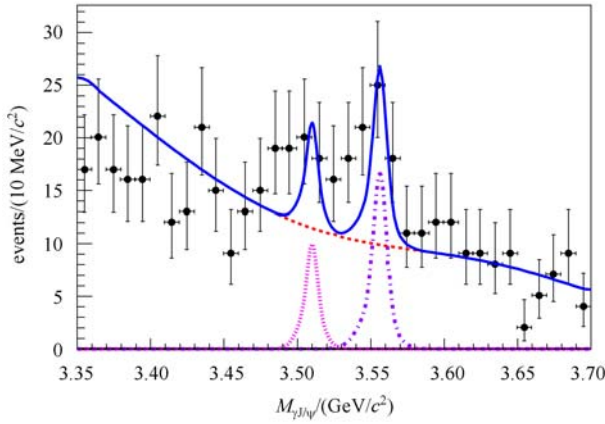


Fig. 3. Result of the simultaneous fit to $M_{\gamma J/\psi}$ distributions for all CME data sets assuming that the signals are from decays of the $Y(4260)$. The blue solid line is the total fit result. The χ_{cJ} signals are shown as dotted line and dash-dotted line, respectively, and the background is shown as the red dashed line.

which all the channels have partially in common. The common sources of systematics include the luminosity measurement, reconstruction efficiencies for charged tracks and photons, the vacuum polarization factor, kinematic fit and branching fractions of the decay of the intermediate states. The systematic uncertainty due to the luminosity measurement is estimated to be 1.0% using Bhabha events [14]. The uncertainty related to the track reconstruction efficiency of high-momentum muons is 1.0% per track [27]. The systematic uncertainty related to the photon detection is estimated to be 1.0% per photon [14]. The systematic uncertainty due to 5C-fit is 0.6%, obtained by studying a control sample of $\psi(2S) \rightarrow \eta J/\psi$ decays. The uncertainty related to the branching fractions of χ_{cJ} and J/ψ decays are taken from the PDG [21]. The uncertainty for the vacuum polarization factor is 0.5% [26].

The other systematic uncertainties arising from the χ_{cJ} mass resolution, the shift of the χ_{cJ} reconstructed mass, the MC model, the shape of the background, the radiative correction factor and the fit range at different CME points are discussed below.

The $\psi(2S) \rightarrow \gamma \chi_{cJ}$ channel is employed as a control sample to extract the differences on the mass resolution of the χ_{cJ} signal by fitting the $M_{\gamma J/\psi}$ spectrum. The differences in the mass resolutions between data and MC are found to be 0.02%, 0.01%, 0.2% for χ_{cJ} ($J=0, 1, 2$). A similar fit is performed, in which the signal shapes are smeared to compensate for the mass resolution difference, and the differences on the yields of χ_{cJ} signal are taken as the systematic uncertainties due to the mass resolution.

An alternative fit is performed shifting the mean of χ_{cJ} signal shapes by one standard deviation of the PDG

values, and the deviations of the signal yields to the nominal values are taken as the systematic uncertainties due to the uncertainties of the signal line shapes.

The detection efficiency is evaluated using MC samples based on the E1 transition assumption [19] for $Y(4260) \rightarrow \gamma \chi_{cJ}$. Another set of MC samples is generated where the $Y(4260) \rightarrow \gamma \chi_{cJ}$ decay is modeled using a phase space distribution, and the differences of the detector efficiencies between the two sets of MC samples are treated as systematic uncertainties from the MC model.

To estimate the systematic uncertainty related to the background shape, a control sample is selected from the data by requiring a $\mu^+ \mu^-$ pair and at least one photon. An alternative background shape is then extracted by re-weighting the $\gamma \mu^+ \mu^-$ invariant mass spectrum of the control sample, where the weights are the efficiency ratio of $e^+ e^- \rightarrow (n\gamma) \mu^+ \mu^-$ MC simulated events surviving the signal selection criteria to the same selection criteria for the control sample. A fit with the alternative background shape is performed, and the differences between the yields of χ_{cJ} signal to the nominal ones are taken as the systematic uncertainties due to the shape of the background.

The possible distortions of the $Y(4260)$ line shape due to interference effects with nearby resonances could introduce uncertainties in the radiative correction factor $\epsilon \times (1 + \delta^r)$. To estimate the related systematic uncertainties, we instead assume that $e^+ e^- \rightarrow \gamma \chi_{cJ}$ are produced via $\psi(4040)$ decays at $\sqrt{s}=4.009$ GeV, $\psi(4160)$ decays at $\sqrt{s}=4.229$ and 4.260 GeV, and $\psi(4415)$ decays at $\sqrt{s}=4.360$ GeV. The variations in the factor $\epsilon \times (1 + \delta^r)$ are taken as the systematic uncertainties due to the radiative correction factor.

A series of similar fits are performed in different ranges of the $M_{\gamma J/\psi}$ distribution, and the largest differences on the signal yields to the nominal values are taken as systematic uncertainties.

All the systematic uncertainties from the different sources are summarized in Table 3. The total systematic uncertainties are calculated as the quadratic sum of all individual terms.

8 Summary

Using data samples collected at CME of $\sqrt{s}=4.009, 4.230, 4.260$, and 4.360 GeV with the BESIII detector, we perform a search for $e^+ e^- \rightarrow \gamma \chi_{cJ}$ ($J=0, 1, 2$) with the subsequent decay $\chi_{cJ} \rightarrow \gamma J/\psi$ and $J/\psi \rightarrow \mu^+ \mu^-$. We find evidence for the processes $e^+ e^- \rightarrow \gamma \chi_{c1}$ and $e^+ e^- \rightarrow \gamma \chi_{c2}$ with statistical significances of 3.0σ and 3.4σ , respectively. No evidence of $e^+ e^- \rightarrow \gamma \chi_{c0}$ is observed. The corresponding Born cross sections of $e^+ e^- \rightarrow \gamma \chi_{cJ}$ at different CME are calculated and listed in Table 2. Under the assumption of the absence of χ_{cJ} signals, the upper

limits on the Born cross sections at the 90% C.L. are calculated and listed in Table 2, too. These upper limits on the Born cross section of $e^+e^- \rightarrow \gamma\chi_{cJ}$ are compatible with the theoretical prediction from an NRQCD calcu-

lation [16].

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